

# Thermometer screen intercomparison in De Bilt (the Netherlands) – Part II: Description and modeling of mean temperature differences and extremes

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**ABSTRACT:** Temperatures of nine thermometer screens during a 6-year field experiment in De Bilt (the Netherlands) have been compared. The screens are currently in use throughout the world and comprise the following types: an aspirated Young screen, four naturally ventilated round-shaped multi-plate screens (KNMI, Vaisala, Young, Socrima), a slightly aspirated version of the KNMI screen, a synthetic Stevenson screen (both aspirated and naturally ventilated) and a naturally ventilated wooden Stevenson screen. The multi-plate KNMI screen served as a reference. For the daily minimum, maximum, and mean air temperatures ( $T_n$ ,  $T_x$ ,  $T_{\text{mean}}$ ), the absolute seasonal mean differences with the reference were almost all  $\leq 0.1$  °C. An exception is the aspirated Young screen for which differences in mean  $T_x$  in summer (JJA) are most notable and amount to  $-0.43$  °C. The differences of the individual  $T_n$ ,  $T_x$ ,  $T_{\text{mean}}$  values may be much larger than their seasonal averages. For the aspirated Young and the naturally ventilated Stevenson screens they are most notable, where the Young is generally cooler ( $T_x$  and  $T_{\text{mean}}$ ) and the Stevensons warmer than the reference. The absolute temperature differences between the screens and the reference are shown to increase with decreasing cloud cover and windspeed. Furthermore, using the original 15-s samples it is shown that random variations cause fast-responding screens to have more extreme  $T_n$ ,  $T_x$  values than slow-responding screens. For a supposed transition of the natural ventilated synthetic Stevenson screen to the reference screen, transfer functions are successfully developed for the 10-min temperature data and the daily  $T_n$  and  $T_x$  data. It is argued that the explained variance of the models could have further improved when high-accuracy (especially in the range 0–3 m/s) windspeed measurements were available (at screen level and position). Copyright © 2007 Royal Meteorological Society

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## 1. Introduction

Increasing interest in air temperature trends and extremes (Klein Tank and Können, 2003; Alexander *et al.*, 2006) necessitates an adequate assessment of inhomogeneities in temperature measurements on a daily to sub-daily level. A potentially important source of these inhomogeneities is the change in thermometer (or radiation) screen design that currently takes place in many countries around the World. Because of ongoing automation, smaller and economically more attractive screens replace traditional screens like the Stevenson screen. An unwanted spin-off of these new, often automated, systems may be the introduction of inhomogeneities in the temperature time series. These inhomogeneities may be of the same order of magnitude as the long-term temperature trends (e.g. Quayle *et al.*, 1991) and may therefore seriously restrict the usefulness of the data.

Screen changes will most likely result in measurements that are closer to the real air temperature, which were defined in Part I as ‘the temperature of the air at the

position of the sensor if no measurement equipment would be installed’ (Van der Meulen and Brandsma, 2007). It is well known that the differences between the temperatures measured in a screen and the real air temperature are largest (up to several degrees Celsius) for clear-sky and windless conditions and that they reduce to zero for overcast and windy conditions (see, e.g. Brandsma, 2004; Parker, 1994, 2004). It logically follows that the magnitude of the inhomogeneities in time series of air temperature, resulting from changes in thermometer screens, strongly depends on the weather conditions. To homogenize the series afterwards on a daily or sub-daily level, weather-dependent corrections have to be applied. However, the development of homogenization methods for that purpose is just starting (Brandsma and Können, 2006; Della-Marta and Wanner, 2006).

Instead of statistically homogenizing temperature time series afterwards, it may be better to derive and apply transfer functions that transform the temperatures of one screen into those of the other. A logical way to proceed may then be to perform a field intercomparison of the old and the new screen. The data of the intercomparison can then be used to develop the transfer function. It is

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not yet clear how long such an intercomparison should take, but sometimes a period of about 1–2 years is suggested (WMO, 2003). The transfer functions should be derived from an analysis of the inter-screen temperature differences in combination with simultaneously measured other weather variables of interest. This labor-intensive work should only be done for a selected set of climate monitoring stations.

Thermometer screen intercomparisons have been reviewed by several authors (Mawley, 1897; Sparks, 1972; Parker, 1994; Barnett *et al.*, 1998). Most of the reported studies focus on the thermometer screens which were in use in the past and the emphasis is on mean monthly temperature differences between screens. National studies of the last 15 years focus on the thermometer screens currently in use (e.g. Andersson and Mattisson, 1991; Lefebvre, 1998; Hatton, 2002; Larre and Hegg, 2002). In these latter papers attempts sometimes have been made to study the temperature differences between screens for varying weather situations. The magnitude of the inter-screen temperature differences in the above studies largely depends on the screens compared. Comparison of Stevenson screens (or similar ones) with outdated screens may reveal mean monthly differences in daily, maximum, and minimum temperatures between a few tenths of degree Celsius in winter to about one degree in summer. On the other hand, comparison of Stevenson screens with modern small round-plated screens shows much smaller mean temperature differences, ranging between zero and about two tenths of a degree Celsius.

Although there have been a number of intercomparisons of thermometer screens around the world, only a few have been published in the international literature and little is quantitatively known about the dependence of the mutual differences between screens on other weather variables. Also a balanced treatment of short-term effects (e.g. for a particular day with specific weather conditions) and mean effects (averages for a large number of days) is lacking. In this study, we present the results of a 6-year field experiment in De Bilt (the Netherlands) that originally compared ten screens that are currently in use around the world. In Part I (Van der Meulen and Brandsma, 2007), we focused on understanding inter-screen temperature differences for typical days. As a result of the calibration in that study one of the screens had to be omitted, and the present analysis is therefore restricted to nine screens. We focus here on the mean temperature differences between the screens and extremes in relation to other weather variables. For the screens and climate conditions considered, the study enables the climate-research community to assess whether modern-day changes of thermometer shelters have caused (or will cause) inhomogeneities in average and extreme temperatures. If there are no significant changes, other countries may save the effort and costs of making comparisons themselves. For a case where the changes are significant, we derive transfer functions that can be used to transform the temperature data of one screen into the other.

In the remainder of the paper, Section 2 first describes the data and methods used. Section 3 presents the results of the field experiment. Section 4 concludes the paper with a discussion and conclusions.

## 2. Data and Methods

### 2.1. Data

The data used in the screen comparison were collected during a 6-year field experiment (9 January, 1989–11 February, 1995) at the testing-site of KNMI in De Bilt (52°06'N, 05°11'E), situated in the center of the Netherlands. A detailed presentation of the setup of the experiment and the calibration of the temperature data is presented in Part I (Van der Meulen and Brandsma, 2007). This section only presents the most essential information.

All temperature measurements were performed at a height of 1.5 m above a flat terrain with short cut-grass cover and sufficiently far removed from major obstacles. Nine screens are compared in this study. Details of the screens are presented in Table I. The KNMI multi-plate screen (Knmi.ref) operated during the whole 6-year period and is taken as the reference screen. The other screens operated for periods of at least two years. Three major types of screens can be distinguished from Table I: Stevenson screens, multi-plate screens, and the Young aspirated screen. The latter screen has often been suggested as an ideal reference screen.

The operational measurement uncertainty of the sensors is 0.1 °C. Owing to the calibration described in part I, we were able to obtain inter-sensor accuracies of about 0.03 °C. Temperature is sampled every 15 s. Unless stated otherwise, we use the 10-min mean temperature values of these samples. Besides temperature, we use the following operationally measured elements at the KNMI-terrain: horizontal wind speed  $u$  at 10 m until 26 June 1993 and thereafter at 20 m, global radiation or total short-wave radiation received at the surface  $K\downarrow$ , cloud cover  $N$ , and relative humidity  $rh$ .

The terrain roughness is large, especially for directions between southeast and north, where 15–30 m high trees are found at distances ranging between 80 and 220 m from the measurement site. Operational wind measurements at the terrain of KNMI are therefore taken about 200 m east from the experimental site and the other operationally measured elements. As a result, the wind measurements may not always be representative of the actual wind speed at screen height near the experimental site.

### 2.2. Methodology

For most climatological applications the main interest is in the differences in daily minimum ( $T_n$ ), daily maximum ( $T_x$ ), and daily mean ( $T_{\text{mean}}$ ) temperature. These differences will be presented first in Section 3. For each screen, we calculated the temperature differences  $\Delta T$  (screen – Knmi.ref) for  $T_n$ ,  $T_x$ , and  $T_{\text{mean}}$  for days with complete data for both the screen and Knmi.ref.  $T_n$

Table I. Details of screens and sensors. The Stevenson screens are of KNMI design.

Screen	Abbreviation	Start date	End date	Diameter (m)	Ventilation	Sensor
KNMI multi-plate	Knmi.ref	89/01/09	95/02/01	0.30	Natural	Pt500
KNMI multi-plate aspirated	Knmi.asp	90/12/13	93/02/20	0.30	Aspirated (1 dm <sup>3</sup> /min)	Pt500
Vaisala multi-plate DTR11	Vaisala	89/01/09	93/02/20	0.30	Natural	Pt500
Young Gill multi-plate 41002	Young	89/01/09	93/02/20	0.12	Natural	Pt500
Young aspirated type II 43408	Young.aspII	92/08/18	95/02/01	0.15/0.025	Aspirated (0.1 dm <sup>3</sup> /s)	Pt1000 <sup>1</sup>
Socrima multi-plate BMO 1167A	Socrima	91/03/08	95/02/01	0.20	Natural	Pt500
Stevenson PVC	Stev.pvc	89/01/09	91/03/06	0.70	Natural	Pt500
Stevenson PVC aspirated	Stev.pvc.asp	91/03/07	95/02/01	0.70	Aspirated 1 dm <sup>3</sup> /s)	Pt500
Stevenson wood	Stev.wood	89/01/09	93/02/20	0.70	Natural	Pt500

<sup>1</sup> For the Young aspirated screen, the sensor is an integral part of the measuring device. The first value for the diameter of these screens refers to the diameter of the radiation shield, the second to the tube that houses the sensor.

and  $T_x$  are calculated for each screen separately and for each day as the minimum and maximum, respectively, of the 144 non-sliding 10-min temperatures.  $T_{\text{mean}}$  is calculated as the mean of these 10-min temperatures. A day runs from 0–24 UTC.

After the presentation of results for  $T_n$ ,  $T_x$ , and  $T_{\text{mean}}$ , the differences in the daily cycle for each season were studied for four combinations of windspeed and cloudiness: I:  $u \leq 3.5 \text{ m/s} \cap N \leq 4/8$ ; II:  $u > 3.5 \text{ m/s} \cap N \leq 4/8$ ; III:  $u \leq 3.5 \text{ m/s} \cap N \geq 5/8$ ; and IV:  $u > 3.5 \text{ m/s} \cap N \geq 5/8$ , where  $N$  equals the fraction of cloud cover (originally in octas).

A fact that is often overlooked is the effect of random temperature variations on the calculation of  $T_n$  and  $T_x$ . The original temperature samples are influenced by turbulent eddies with a wide range of space and length scales. Those eddies are advected along the screens causing deviations from the mean temperature. It is obvious that fast-responding screens follow better these turbulence driven temperature deviations than slower ones. In practice,  $T_n$  and  $T_x$  are operationally calculated from sliding  $L$ -min averages of the temperature samples. The turbulence driven temperature deviations cause the

$L$ -min averages of fast-responding screens to fluctuate more around the mean temperature than those of slow-responding screens. As a result, fast-responding screens have a higher probability of obtaining more extreme  $T_n$  and  $T_x$  than slow-responding screens (disregarding other causes of temperature differences, like radiation errors). This effect was studied by comparing the screens for several lengths of the averaging interval using the original 15-s data.

Parallel measurements like those presented here can be used in climatology for two main reasons: (1) to assess the magnitude of the bias for the transition from an old screen to a new screen, and (2) to derive a transfer function that transforms the temperatures measured with the old screen into a series measured with the new screen, or the other way around. Here we derive such a transfer function. In Part I, we demonstrated the importance of lag-time differences, windspeed, and cloudiness in explaining temperature differences between screens. Here we use that information and the results in the present paper for the derivation of the transfer function. Stev.pvc and Knmi.ref are used here as an example.

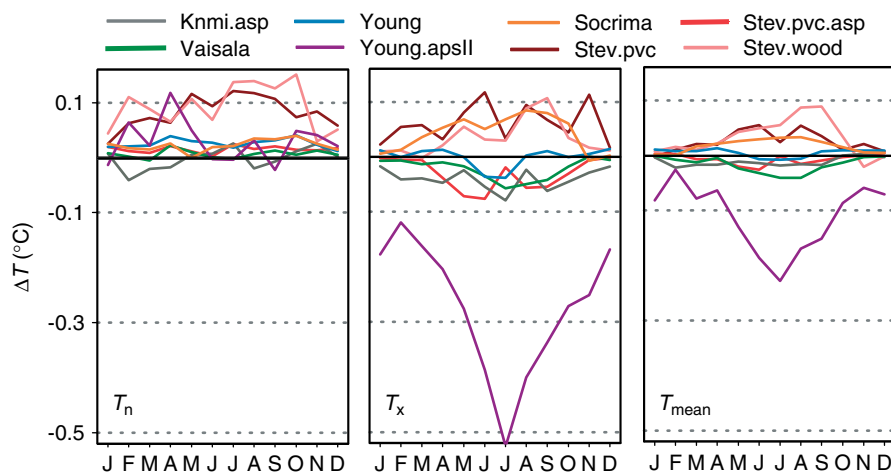


Figure 1. Monthly mean temperature differences  $\Delta T$  (screen – Knmi.ref) for daily minimum temperature  $T_n$ , maximum temperature  $T_x$ , and mean temperature  $T_{\text{mean}}$ . The screens and their period of overlap with Knmi.ref are defined in Table I.

Table II. Mean temperature differences  $\Delta T$  (screen – Knmi.ref) for winter (DJF), spring (MAM), summer (JJA), and autumn (SON) for daily minimum temperature  $T_n$ , maximum temperature  $T_x$ , and mean temperature  $T_{mean}$ . The screens and their period of overlap with Knmi.ref are defined in Table I. The values in brackets give the standard error.

	$\Delta T_n$ (°C)				$\Delta T_x$ (°C)				$\Delta T_{mean}$ (°C)			
	DJF	MAM	JJA	SON	DJF	MAM	JJA	SON	DJF	MAM	JJA	SON
Knmi.asp	-0.009 (0.007)	-0.010 (0.008)	0.004 (0.006)	0.006 (0.008)	-0.024 (0.004)	-0.036 (0.006)	-0.052 (0.006)	-0.049 (0.006)	-0.009 (0.002)	-0.015 (0.002)	-0.016 (0.002)	-0.008 (0.002)
Vaisala	0.005 (0.002)	0.010 (0.003)	0.002 (0.003)	0.010 (0.003)	-0.007 (0.002)	-0.013 (0.004)	-0.047 (0.005)	-0.021 (0.005)	-0.003 (0.001)	-0.014 (0.001)	-0.038 (0.001)	-0.013 (0.001)
Young	0.018 (0.003)	0.031 (0.005)	0.024 (0.005)	0.032 (0.005)	0.009 (0.004)	0.007 (0.005)	-0.024 (0.006)	0.005 (0.005)	0.009 (0.001)	0.009 (0.002)	-0.007 (0.002)	0.009 (0.001)
Young.aspII	0.019 (0.015)	0.065 (0.016)	0.010 (0.015)	0.022 (0.012)	-0.159 (0.015)	-0.220 (0.017)	-0.433 (0.021)	-0.286 (0.019)	-0.064 (0.009)	-0.093 (0.009)	-0.192 (0.007)	-0.099 (0.008)
Socrima	0.019 (0.003)	0.013 (0.004)	0.025 (0.004)	0.033 (0.004)	0.006 (0.004)	0.054 (0.004)	0.068 (0.005)	0.048 (0.007)	0.002 (0.001)	0.022 (0.002)	0.031 (0.002)	0.016 (0.002)
Stev.pvc	0.046 (0.009)	0.085 (0.010)	0.111 (0.012)	0.088 (0.011)	0.032 (0.013)	0.057 (0.011)	0.081 (0.013)	0.070 (0.014)	0.004 (0.003)	0.030 (0.004)	0.044 (0.005)	0.022 (0.004)
Stev.pvc.asp	0.017 (0.003)	0.014 (0.004)	0.014 (0.005)	0.016 (0.004)	-0.003 (0.004)	-0.042 (0.005)	-0.051 (0.011)	-0.032 (0.005)	0.003 (0.001)	-0.012 (0.002)	-0.015 (0.004)	-0.004 (0.001)
Stev.wood	0.065 (0.012)	0.086 (0.012)	0.115 (0.013)	0.101 (0.013)	0.011 (0.008)	0.025 (0.009)	0.049 (0.010)	0.053 (0.012)	0.006 (0.003)	0.025 (0.003)	0.063 (0.004)	0.034 (0.006)

\* Values in italics are not significantly different from zero ( $2 \times se$ ).

### 3. Results

#### 3.1. Mean and extreme differences in daily $T_n$ , $T_x$ , and $T_{mean}$

Figure 1 shows the monthly mean  $\Delta T$  values. The most striking feature is the anomalous behaviour of  $\Delta T_x$  for Young.aspII (also reflected in  $T_{mean}$ ), with the largest deviation from Knmi.ref in July when  $T_x$  of Young.aspII is on average 0.5 °C lower than  $T_x$  of Knmi.ref. The temperatures for the multi-plate screens are all close to Knmi.ref. Socrima, however, has a tendency to be a little warmer than the other multi-plate screens, which may be related to reduced natural ventilation within this screen compared to the other multi-plate screens. The natural ventilated Stevenson screens (Stev.pvc, Stev.wood) are warmer than Knmi.ref both during day and night. The monthly mean  $\Delta T_n$  and  $\Delta T_x$  are largest in the summer half year, but mostly do not exceed 0.1 °C. The results for Stev.pvc.asp show that the effect of aspiration is that the Stevenson screen more closely resembles the multi-plate screens compared to the unventilated version of the screen (Stev.pvc).

The  $\Delta T$  values for each screen are presented in Table II (seasonal average values) and Table III (annual average values). From the 96 average  $\Delta T$  values in Table II, 80 are significantly different from zero (because of serial correlation in the difference series within a season, the standard errors may go up to 1.5 times larger than those presented). However, only 37 of these values are larger than the inter-sensor accuracy (0.03 °C). As may already be inferred from Figure 1, the largest  $\Delta T$  values occur for Young.aspII, Stev.pvc, and Stev.wood. The annual average  $\Delta T$  values in Table III are small for these three screens but not always negligible. For  $\Delta T_n$  they range between 0.029 °C (Young.aspII) and 0.092 °C (Stev.wood); for  $\Delta T_x$  between -0.275 °C (Young.aspII) and 0.060 °C (Stev.pvc); and for  $\Delta T_{mean}$  between -0.112 °C (Young.aspII) and 0.032 °C (Stev.wood).

Additional information about the distribution of the individual daily  $\Delta T$  values can be obtained from the

Table III. Annual mean temperature differences  $\Delta T$  (screen – Knmi.ref) for daily minimum temperature  $T_n$ , maximum temperature  $T_x$ , and mean temperature  $T_{mean}$ . The screens and their period of overlap with Knmi.ref are defined in Table I. The values in brackets give the standard error.

	$\Delta T_n$ (°C)	$\Delta T_x$ (°C)	$\Delta T_{mean}$ (°C)
Knmi.asp	-0.002 (0.004)	-0.040 (0.003)	-0.012 (0.001)
Vaisala	0.007 (0.001)	-0.022 (0.002)	-0.017 (0.000)
Young	0.026 (0.002)	-0.001 (0.003)	0.005 (0.001)
Young.aspII	0.029 (0.007)	-0.275 (0.009)	-0.112 (0.004)
Socrima	0.023 (0.002)	0.044 (0.003)	0.018 (0.001)
Stev.pvc	0.082 (0.005)	0.060 (0.006)	0.025 (0.002)
Stev.pvc.asp	0.015 (0.002)	-0.032 (0.003)	-0.007 (0.001)
Stev.wood	0.092 (0.006)	0.035 (0.005)	0.032 (0.002)

\* Values in italics are not significantly different from zero ( $2 \times se$ ).

boxplots in Figure 2. Consider for example the boxplot of  $\Delta T_n$  of Stev.wood for autumn (SON). This boxplot shows that the  $\Delta T_n$  distribution is positively skewed with the 90<sup>th</sup> percentile equal to 0.33 °C (in 10% of the days  $\Delta T_n$  is larger than this number). A result of the positive skewness in this case is that the average  $\Delta T_n$  of 0.10 °C (Table II) is 0.04 °C larger than the median (50<sup>th</sup> percentile) of 0.06 °C shown in the boxplot. The figure shows that most of the  $\Delta T$  distributions of Stev.pvc and Stev.wood exhibit this positive skewness. Figure 2 also shows that the distribution of  $\Delta T$  values for  $T_{\text{mean}}$  is much narrower than that for  $T_n$ ,  $T_x$ . For (almost) identical screens, the boxplots may also provide information about the spread of the  $\Delta T$  values that may be expected because of pure random variation alone. In this experiment, Vaisala is more or less identical to Knmi.ref and the corresponding statistics of the  $\Delta T$  values between these screens could be considered as a measure for natural background noise.

### 3.2. Diurnal temperature cycle and the effect of wind and cloudiness

For each screen, the annual mean diurnal cycle of the temperature differences  $\Delta T$  (screen – Knmi.ref) was

calculated with an hourly resolution. Figure 3 shows these cycles for all screens together with the mean annual diurnal temperature cycle of Knmi.ref. For Stev.pvc and Stev.wood, it appears that the largest absolute values of  $\Delta T$  are just after sunrise and just after sunset. Because of their relatively large response times, these screens have more difficulties in following the steep temperature changes at those times than the other screens (see also Part I). For Young.aspII, the largest  $\Delta T$  is around the time of maximum solar radiation. Figure 4 shows the same diurnal cycle differences but now for the four combinations of  $u$  and  $N$  defined in Section 2.2 for both winter (DJF) and summer (JJA). The figure clearly shows that the temperature differences between screens depend on  $u$  and  $N$ . Small values of  $u$  and  $N$  (category I) result in large mutual differences in the diurnal cycles while large values (category IV) minimize those differences. It is interesting to see that the large absolute values of  $\Delta T$  for Stev.pvc and Stev.wood, just after sunrise and just after sunset, almost disappear for  $u > 3.5$  m/s. It is also noteworthy that for large values of  $u$  and  $N$  (category IV) the daytime bias of Young.aspII remains relatively large (about  $-0.1$  °C in winter and  $-0.2$  °C in summer).

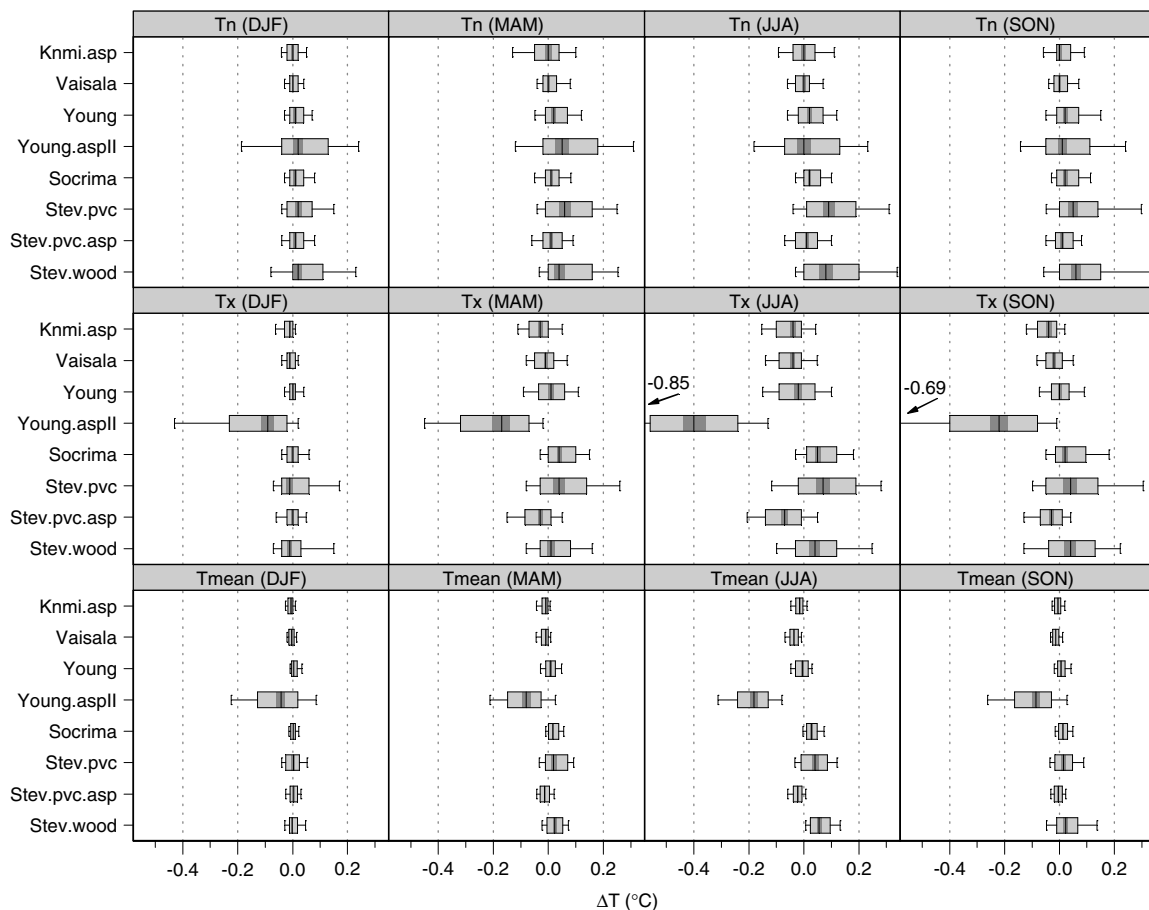


Figure 2. Boxplots of the individual temperature differences  $\Delta T$  (screen – Knmi.ref) for winter (DJF), spring (MAM), summer (JJA), and autumn (SON) for daily minimum temperature  $T_n$ , maximum temperature  $T_x$ , and mean temperature  $T_{\text{mean}}$ . The screens and their period of overlap with Knmi.ref are defined in Table I. The left and right limits of the box represent the 25th/75th percentiles (quartiles); the vertical line within the box represents the 50th percentile (median) with 95% confidence interval (dark gray); and the whiskers mark the 10th/90th percentiles. Two of the Young.aspII whiskers are outside the horizontal scale; the corresponding values are shown in the figure with arrows pointing to the location of the whiskers.

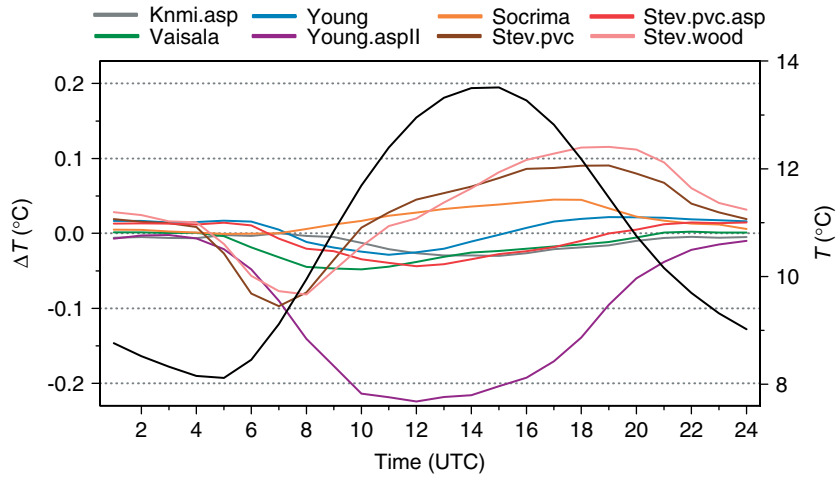


Figure 3. Annual mean diurnal cycles of the hourly temperature differences  $\Delta T$  (screen – Knmi.ref) of the eight screens (left axis) and of the hourly temperatures of Knmi.ref (black line, right axis). The screens and their period of overlap with Knmi.ref are defined in Table I.

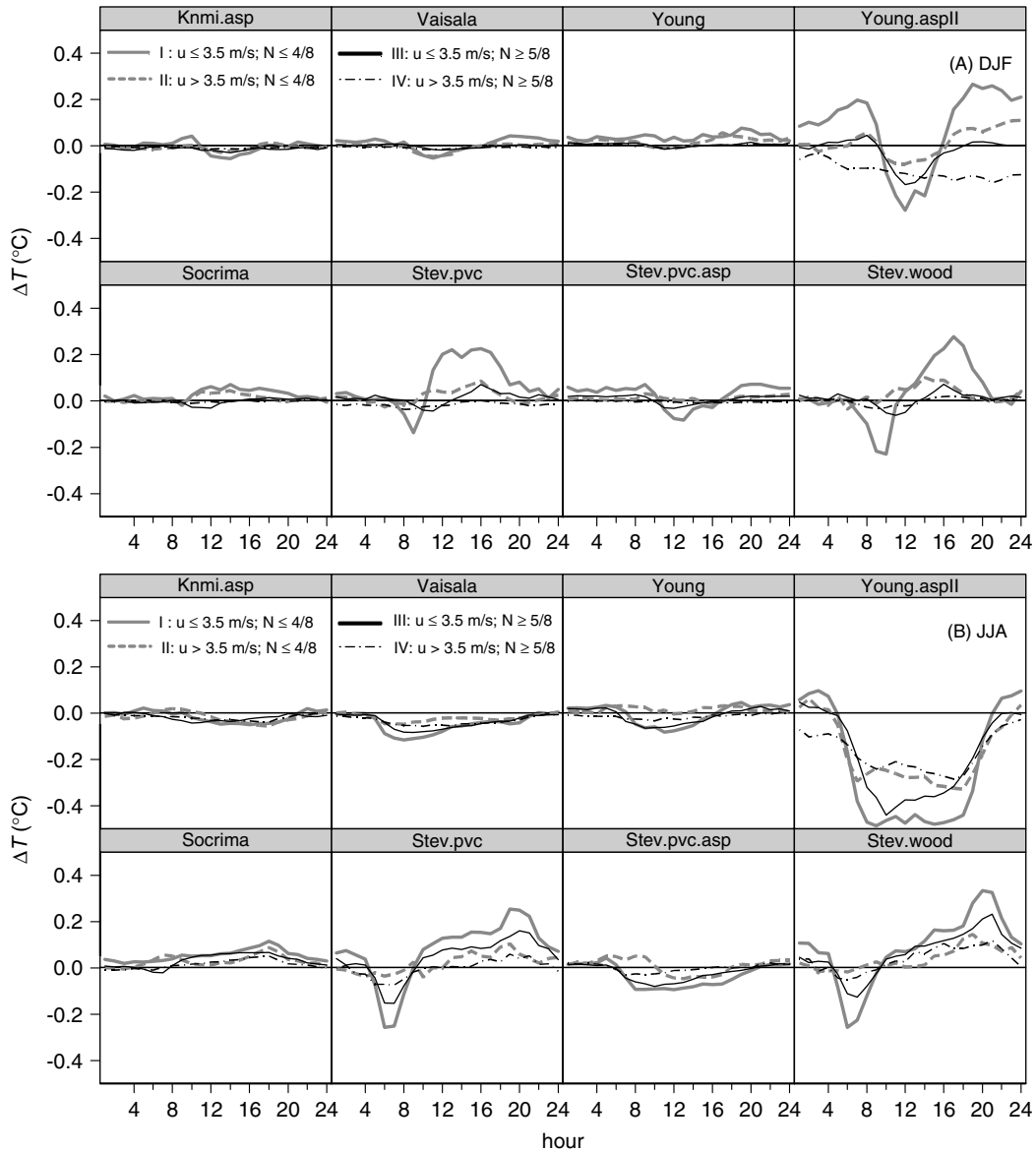


Figure 4. Mean diurnal cycles of the hourly temperature differences  $\Delta T$  (screen – Knmi.ref) for four combinations of windspeed ( $u$ ) and fraction of cloudiness ( $N$ ). The upper two rows (A) apply to the three winter months (DJF) and the lower two rows (B) to the three summer months (JJA). The screens and their period of overlap with Knmi.ref are defined in Table I.

3.3. Effect of averaging interval length  $L$  on  $T_n$  and  $T_x$ . The effect of the averaging interval length  $L$  on  $T_n$  and  $T_x$  is studied by calculating  $T_n$  and  $T_x$  for  $L = 0.25, 1, 2, \dots, 15$  min, where  $L = 0.25$  min corresponds to the original 15-s temperature samples (no averaging).  $T_n$  and  $T_x$  are selected as the smallest and largest values, respectively, from the sliding  $L$ -min temperature averages in a day ( $24 \times 60 \times 4 = 5760$  values). For the presentation of the results, we restrict ourselves to two thirty-day summer periods, without missing values, in 1989 (July 1–30) and in 1994 (June 17–July 16). For each  $L$ ,  $T_n$  and  $T_x$  are calculated as the average of the 30 daily values. To facilitate the comparison between the screens, we subtracted for each screen the mean  $T_n$  and  $T_x$  values for all  $L$ . The resulting anomalies are denoted as  $T_n^*$  and  $T_x^*$ .

Figure 5 shows the results for July 1989 for both  $T_n^*$  (left) and  $T_x^*$  (right). The figure clearly shows that the modern type round multi-plated screens are more sensitive to changes in  $L$  than the Stevenson screens and that these effects are much larger for  $T_x^*$  than for  $T_n^*$ . For instance, shifting from  $L = 10$  min to  $L = 1$  min results in an increase in  $T_x$  of Knmi.ref of  $0.33^\circ\text{C}$ , whereas for Stev.pvc the increase amounts to only  $0.13^\circ\text{C}$ . For  $T_n$ , the corresponding values are a decrease in  $T_n$  of Knmi.ref of  $0.10^\circ\text{C}$  and a decrease of Stev.pvc of  $0.03^\circ\text{C}$ .

Figure 6 shows the results in 1994 for the period between June 17 and July 16. Compared to the July 1989 period, the absolute values for Knmi.ref are somewhat smaller here. Noteworthy are the large (absolute) values for Young.aspII and (to a smaller extent) of Stev.pvc.asp

in comparison to Knmi.ref. In January (not shown) the effects of changing  $L$  are about a factor 3–4 smaller than in July. Note also the relatively small absolute values for Socrima compared to Knmi.ref.

The implication of the above is that fast-responding screens have by nature more extreme  $T_x$  and  $T_n$  than slow-responding screens and they are more sensitive to changes in  $L$ . This effect may result in differences of several tenths of  $^\circ\text{C}$  between the  $T_x$  and  $T_n$  of those screens. In practice, this effect has to be considered in connection with other effects (e.g. the larger artificial heating of natural ventilated Stevenson screens compared to modern screens).

### 3.4. Transfer function for a transition of Stev.pvc to Knmi.ref

A supposed transition of Stev.pvc to the faster-responding Knmi.ref is used here to demonstrate the construction of transfer functions. Stev.pvc is selected because its temperatures show significant differences with those of Knmi.ref and because it was in use for an almost uninterrupted period of nearly two years without major problems becoming visible in the calibration (see also Part I). The period of simultaneous measurements is used to derive a transfer function that enables a user to transform the new measurements into the old ones or the other way around. We consider two transfer functions: (1) for the 10-min temperatures, and (2) for  $T_n$  and  $T_x$ . In both cases, we first explored the relationships between the predictant and all possible predictors for the

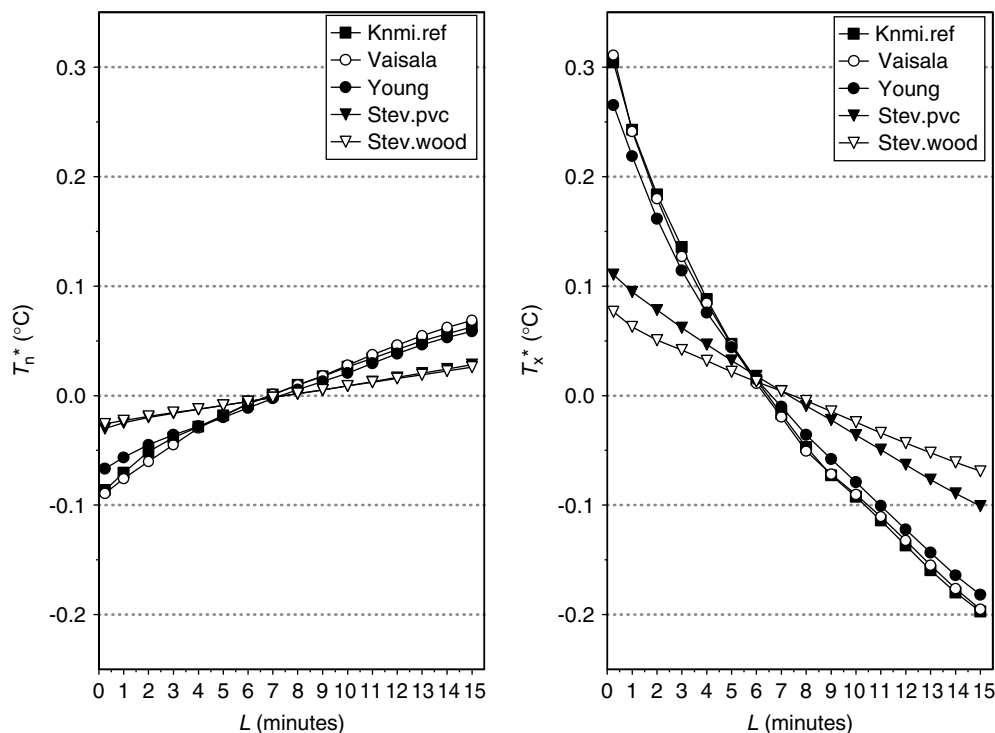


Figure 5. Mean daily  $T_n^*$  (left) and  $T_x^*$  (right) as a function of the averaging interval for July 1989. For each screen, the mean  $T_n$  and  $T_x$  values for all  $L$  are subtracted to obtain  $T_n^*$  and  $T_x^*$ . Each daily  $T_n$  and  $T_x$  is calculated from 5760 overlapping (15 s) intervals of length  $L = 0.25, 1, 2, \dots, 15$  min.

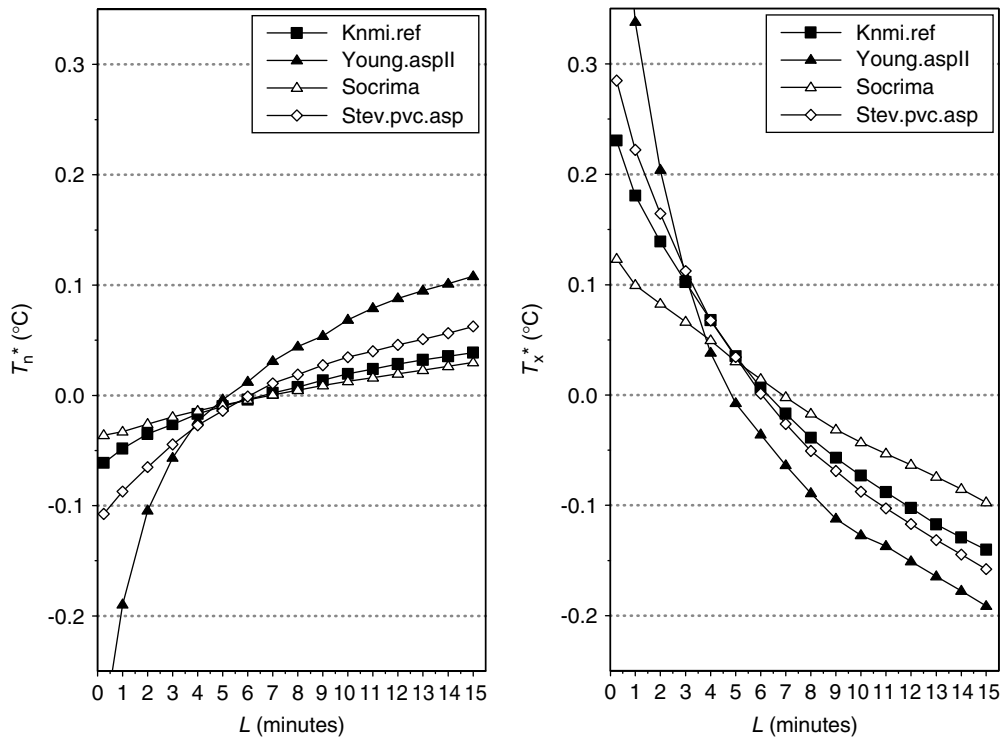


Figure 6. As in Figure 5, but showing the period ranging from June 17–July 16, 1994. The  $L = 0.25$  min values for Young.asplI fall out of the paper and equal  $-0.30^\circ\text{C}$  in the left plot and  $0.57^\circ\text{C}$  in the right plot.

presence of non-linear effects and interaction using non-parametric loess functions (Cleveland and Devlin, 1988). After fitting the equations, we examined the relationships between the error series and the predictors using the same procedure.

### 3.4.1. Model for 10-min data

For the 10-min data we propose the following equation:

$$\Delta T_i = a_0 + a_1 \Delta T_{i-1} + a_2 \frac{(dT_{ref}/dt)_i}{(u_i + 1)} + a_3 \frac{K \downarrow_i / 100}{(u_i + 1)} + a_4 \frac{\min(1 - N_i, I_i)}{(u_i + 1)} + \varepsilon_i$$

$$I_i = \begin{cases} 0 & \text{during daytime} \\ 1 & \text{during nighttime} \end{cases} \quad (1)$$

where  $i$  is the counter of the 10-min interval considered ( $i \in [1, 2, \dots, M]$ ) with  $M$  the total number of 10-min intervals,  $\Delta T_i$  equals the temperature difference  $T(\text{Stev.pvc}) - T(\text{Knmi.ref})$  in  $^\circ\text{C}$ , for the  $i$ th 10-min interval,  $dT_{ref}/dt$  is the derivative of the temperature curve of Knmi.ref at time  $t$  which is defined for the  $i$ th 10-min interval as  $T(\text{Knmi.ref})_i - T(\text{Knmi.ref})_{i-1}$ ,  $u_i$  is windspeed for the  $i$ th 10-min interval in m/s,  $K \downarrow_i$  global radiation for the  $i$ th 10-min interval in  $\text{W}/\text{m}^2$  which equals zero during nighttime by definition,  $N_i$  is the cloudiness fraction for the  $i$ th 10-min interval ( $N \in [0, 1/8, 2/8, \dots, 8/8]$ ),  $\varepsilon_i$  is the residual for the  $i$ th 10-min interval in  $^\circ\text{C}$  assumed to be independent, normally distributed and zero mean, and  $a_0, \dots, a_4$  are coefficients to be estimated. The two terms with coefficients  $a_1$  and

$a_2$ , on the right-hand side of Equation (1), account for the slow response of Stev.pvc compared to Knmi.ref. The next two terms account for the fact that for clear-sky conditions, observed inter-screen temperature differences may be larger (in an absolute sense) than for cloudy conditions. The division by  $u_i + 1$  in the terms involving  $T_{ref}$ ,  $K \downarrow$ , and  $N_i$  is needed because windspeed dampens their effects on  $\Delta T$ , the  $+1$  m/s is added to account for the fact that anemometers have a threshold speed and to prevent division by zero. As  $K \downarrow$  may take values as large as  $1000 \text{ W}/\text{m}^2$ , it is divided by 100 to facilitate the intercomparison of the coefficients.

The coefficients of the model in Equation (1) have been estimated for each season using ordinary least-squares regression (OLS). Table IV presents the results. For all seasons, the intercept term  $\hat{a}_0$  is significant but nearly zero. The table further shows that there is a strong contribution of the autoregressive term  $\Delta T_{i-1}$  for each season. The relatively large negative values for  $\hat{a}_2$  demonstrate that Stev.pvc is not able to follow the fast temperature fluctuations of Knmi.ref. The clear-sky effect on  $\Delta T$  is significant and seasonally dependent both during day and night. Note that the effects are opposite (relative warming of Stev.pvc during the day and relative cooling during the night) and their potential magnitudes are about an order of magnitude larger during the day (e.g. with conditions in summer with the maximum  $K \downarrow$  about  $1000 \text{ W}/\text{m}^2$  and  $u = 0$  m/s) than during the night (for conditions with  $N = 0$  and  $u = 0$  m/s).

Table IV shows that the explained variance is large and ranges between 81.5% and 87.1% for spring and autumn, respectively. The results do not improve by applying



Table IV. Estimated regression coefficients ( $\hat{a}_0, \dots, \hat{a}_4$ ), number of observations ( $n$ ), residual standard error (RSE), and the coefficient of determination ( $R^2$ ) for the model in Equation (1) for winter (DJF), spring (MAM), summer (JJA) and autumn (SON) for the 10-min temperature differences  $\Delta T$  (Stev.pvc–Knmi.ref) in °C. The values within brackets are the  $t$ -statistics of the parameter estimates above.

	$\hat{a}_0$	$\hat{a}_1$	$\hat{a}_2$	$\hat{a}_3$	$\hat{a}_4$	$n$	RSE	$R^2$
DJF	−0.0093 (−23.6)	0.5918 (208.1)	−1.0378 (−181.5)	0.1682 (80.4)	−0.0233 (−9.0)	27005	0.051	0.816
MAM	−0.0084 (−11.1)	0.5363 (182.0)	−1.0916 (−189.4)	0.0647 (64.2)	−0.0631 (−17.9)	25084	0.082	0.815
JJA	−0.0113 (−13.1)	0.5660 (200.6)	−1.0626 (−187.7)	0.0613 (61.7)	−0.0424 (−11.3)	25962	0.092	0.828
SON	−0.0079 (−12.9)	0.6133 (239.6)	−1.0410 (−199.9)	0.0831 (64.4)	−0.0393 (−13.1)	23703	0.069	0.871

a power transformation to  $(u + 1)$  in Equation (1). The residual series show some autocorrelation with the first-order autocorrelation coefficient ranging between 0.17 in summer and 0.26 in autumn. As a result, the statistical significance of the coefficients in Table IV will in fact be somewhat smaller than the calculated values.

The large values for the explained variance are partly a result of using observed values of  $\Delta T_{i-1}$  in Equation (1). When we use the simulated  $\Delta \hat{T}_{i-1}$  instead (using the estimated regression coefficients  $\hat{a}_0, \dots, \hat{a}_4$ ), the explained variance decreases to values between 56.9% and 66.4% for winter and autumn, respectively. These values do not further degrade when we use the first part of the data (1989) for model calibration and the second part (1990) for model verification.

To check whether model or data deficiencies may (partly) be responsible for the decrease in explained variance, we repeated the experiment using simulated  $\Delta T$  values. First, we simulated a new series of random normally distributed errors  $\varepsilon_i$  in Equation (1) using the calculated mean and standard deviation of  $\varepsilon_i$ . Second, the simulated  $\varepsilon_i$  series was used to simulate a new  $\Delta T_i$  series using the estimated regression coefficients  $\hat{a}_0, \dots, \hat{a}_4$  in Equation (1), and the measured values of the predictors. Finally, Equation (1) was fitted with the new simulated  $\Delta T_i$  series (and the accompanying  $\Delta T_{i-1}$  series) resulting in a new set of estimated regression coefficients  $\hat{a}_0, \dots, \hat{a}_4$ . The explained variance is, as it should be, almost the same as before and ranges between 80.0% and 86.7%. When again using the simulated  $\Delta \hat{T}_{i-1}$  instead of  $\Delta T_{i-1}$  (using the newly estimated regression coefficients  $\hat{a}_0, \dots, \hat{a}_4$ ), the explained variance now decreases much less than before to values between 70.8% and 79.0% for winter and autumn, respectively. These results indicate that model and/or data deficiencies may indeed (partly) be responsible for the decrease in explained variance. Examination of the relationships between the error series of all possible predictors does not reveal serious interaction or non-linear effects. We therefore suggest that the lack of adequate windspeed data at screen height at the experimental site is partly responsible for the decrease in explained variance (windspeed is strongly represented in Equation (1)).

### 3.4.2. Model for $T_n$ and $T_x$ data

At the times  $t_n$  of the minimum temperature and  $t_x$  of the maximum temperature, the first-order derivatives

of the temperature curve equal zero and the effects of differences in response times of the screens on  $\Delta T$  should be minimal. However, in Section 3.3 we have shown that the faster responding screens have a higher probability to reach more extreme  $T_n$  or  $T_x$  than the slow-responding screens. To account for this, we define a new variable reflecting the sharpness of the dale or peak in the temperature curve at  $t_n$  and  $t_x$  of the reference screen:

$$DPS_d = T_{ref,d,i} - \frac{T_{ref,d,i-1} + T_{ref,d,i+1}}{2} \quad (2)$$

where  $DPS_d$  stands for the Dale or Peak Sharpness of Knmi.ref at  $t_n$  or  $t_x$  at day  $d$  and  $T_{ref,d,i}$  is the temperature of Knmi.ref at day  $d$  in the  $i$ th 10-min time interval corresponding to  $t_n$  or  $t_x$ . The larger the absolute value of  $DPS$ , the greater the discrepancy between Knmi.ref and the slower responding Stev.pvc. For  $T_n$  the following equation is used:

$$\Delta T_{n,d} = b_0 + b_1 \frac{DPS_d}{(u_d + 1)} + b_2 \frac{(1 - N_d)}{(u_d + 1)} + \varepsilon_d \quad (3)$$

where  $u_d$  and  $N_d$  are the 10-min windspeed and cloudiness at  $t_n$  for day  $d$ ,  $\varepsilon_d$  is assumed to be an independent normally distributed zero mean residual, and  $b_0, b_1, b_2$  are coefficients to be estimated. Although Equation (3) is in principle set up for nighttime  $t_n$ , it also works for the scarce daytime  $t_n$  values.

For  $T_x$  the following equation is used:

$$\Delta T_{x,d} = c_0 + c_1 DPS_d + c_2 \frac{K \downarrow_d / 100}{(u_d + 1)} + \varepsilon_d \quad (4)$$

where  $u_d$  and  $K \downarrow_d$  are the windspeed and global radiation at  $t_x$  for day  $d$ ,  $\varepsilon_d$  is assumed to be an independent normally distributed zero mean residual, and  $c_0, c_1, c_2$  are coefficients to be estimated. Inclusion of nighttime cloudiness, for nighttime  $t_x$ , is not significant. Division of  $DPS_i$  by  $(u + 1)$  is not advantageous here. Note that  $t_n$  and  $t_x$  of Knmi.ref and Stev.pvc need not be exactly the same. The weather variables in Equations (3) and (4) are at  $t_n$  and  $t_x$  of Knmi.ref.

The coefficients of the models in Equations (3) and (4) have been estimated for each season using OLS. If coefficients turned out to be not significant at the 5% level, the corresponding predictors were omitted from the

equation, starting with the predictor with the smallest  $t$ -statistic. After the deletion of an insignificant predictor, the remaining coefficients were reestimated. Tables V and VI present the results for  $T_n$  and  $T_x$ , respectively. Instead of the coefficient of determination ( $R^2$ ), the tables give the squared correlation coefficient  $r^2$  between the predictant and the fitted values. The latter coefficient is preferred when models with and without intercept are compared (for models with intercept  $R^2 = r^2$ ).

Tables V and VI show that the intercept term is only significant for  $\Delta T_x$  in spring. The coefficients  $\hat{b}_1$ ,  $\hat{b}_2$ ,  $\hat{c}_1$ ,  $\hat{c}_2$  show some seasonal variation. Note that for  $\Delta T_n$  the contribution of the term with cloudiness is not significant in both winter and autumn. For  $\Delta T_n$ , the  $DPS$  term has a relatively large contribution. For instance, in summer  $DPS$  at  $t_n$  ranges between  $-0.7$  and  $0$  resulting in a maximum contribution of  $+0.48^\circ\text{C}$  (with  $u = 0$  m/s) to  $\Delta T_n$ . This is large compared to the maximum contribution of cloudless skies (with  $u = 0$  m/s) of about  $+0.12^\circ\text{C}$  in summer. In contrast, for  $\Delta T_x$  in summer the clear-sky effect has the largest contribution of about  $+1.5^\circ\text{C}$  ( $K \downarrow = 1000 \text{ W/m}^2$  and  $u = 0$  m/s). Here  $DPS$  at  $t_x$  ranges between  $0.0$  and  $0.9$  resulting in a maximum contribution of  $-0.25^\circ\text{C}$  (with  $u = 0$  m/s) to  $\Delta T_x$ .

The  $r^2$  in Tables V and VI are relatively low and range between  $0.27$  and  $0.53$ . We suggest that the lack of adequate windspeed data at screen height at the experimental site is even more important here than for the 10-min model of Equation (1). Another factor could be the relatively large natural variation of  $T$ , which cannot be explained by a model. To obtain an idea of the magnitude of this factor, we can take the difference between Knmi.ref and Vaisala as a measure for the natural variation (Knmi.ref and Vaisala are almost identical screens).

The standard deviation of  $\Delta T_n$  for these screens varies between  $0.035^\circ\text{C}$  in winter and  $0.056^\circ\text{C}$  in summer. For  $\Delta T_x$ , the corresponding values are  $0.037^\circ\text{C}$  in winter and  $0.077^\circ\text{C}$  in summer. Comparing these values with the residual standard error (RSE) in Tables V and VI shows that there is still enough room for improvement.

#### 4. Discussion and conclusions

In this paper, we compared the temperatures of nine thermometer screens that are currently in use around the world. The comparison was based on measurements taken during a 6-year field experiment in De Bilt (the Netherlands). All screens were situated close to each other on the same experimental site and faced the same weather and environmental conditions. Furthermore, the screens were all supplied with the same type of temperature sensor.

The results indicate that the seasonal differences in  $T_n$ ,  $T_x$ , and  $T_{\text{mean}}$ , with respect to the reference screen Knmi.ref, are small (generally  $\leq 0.1^\circ\text{C}$ ) for most screens. This is in agreement with other studies that intercompared similar screens (e.g. Andersson and Mattisson, 1991; Larre and Hegg, 2002). For the artificially ventilated Young.aspII, the seasonal differences in  $T_n$ ,  $T_x$ , and  $T_{\text{mean}}$  are much larger than for all other screens. The value for  $\Delta T_x$  is largest and varies between  $-0.159^\circ\text{C}$  (DJF) and  $-0.433^\circ\text{C}$  (JJA). Also the annual mean  $\Delta T_{\text{mean}}$  is largest for Youngs.apsII and amounts to  $-0.112^\circ\text{C}$  whereas for all other screens the annual mean absolute  $\Delta T_{\text{mean}}$  is not greater than  $0.032^\circ\text{C}$  (Stev.wood). The latter value may be considered small with respect to the observed global temperature rise of the last 150 years.

Table V. Estimated regression coefficients ( $\hat{b}_0, \dots, \hat{b}_2$ ), number of observations ( $n$ ), residual standard error (RSE), and the squared correlation coefficient ( $r^2$ ) for the model in Equation (3) for winter (DJF), spring (MAM), summer (JJA), and autumn (SON) for the daily minimum temperature differences  $\Delta T_n$  (Stev.pvc–Knmi.ref). The values within brackets are the  $t$ -statistics of the parameter estimates above.

	$\hat{b}_0$	$\hat{b}_1$	$\hat{b}_2$	$n$	RSE	$r^2$
DJF	–	$-1.2223$ ( $-12.0$ )	–	171	0.096	0.381
MAM	–	$-0.8018$ ( $-6.9$ )	$0.0741$ ( $2.0$ )	171	0.113	0.306
JJA	–	$-0.6838$ ( $-6.7$ )	$0.1207$ ( $3.3$ )	173	0.130	0.276
SON	–	$-1.1326$ ( $-17.8$ )	–	163	0.099	0.533

Table VI. Estimated regression coefficients ( $\hat{c}_0, \dots, \hat{c}_2$ ), number of observations ( $n$ ), RSE, and the squared correlation coefficient ( $r^2$ ) for the model in Equation (4) for winter (DJF), spring (MAM), summer (JJA), and autumn (SON) for the daily maximum temperature differences  $\Delta T_x$  (Stev.pvc–Knmi.ref). The values within brackets are the  $t$ -statistics of the parameter estimates above.

	$\hat{c}_0$	$\hat{c}_1$	$\hat{c}_2$	$n$	RSE	$r^2$
DJF	–	$-0.4458$ ( $-5.9$ )	$0.5722$ ( $12.2$ )	171	0.126	0.459
MAM	$0.0584$ ( $2.6$ )	$-0.3566$ ( $-6.6$ )	$0.1128$ ( $5.5$ )	170	0.128	0.271
JJA	–	$-0.2826$ ( $-7.7$ )	$0.1531$ ( $12.4$ )	173	0.134	0.363
SON	–	$-0.3648$ ( $-6.9$ )	$0.2695$ ( $11.1$ )	163	0.144	0.349

When looking at extremes, the daily values of  $\Delta T_n$ ,  $\Delta T_x$ , and  $\Delta T_{\text{mean}}$  are important. The results show that these values may be much larger than the seasonal averages and that the studied screens can be roughly divided into three major groups: (1) Knmi.asp, Vaisala, Young, Socrima, and Stev.pvc.asp; (2) Stev.pvc and Stev.wood; and (3) Young.aspII. The screens in the first group behave more or less the same as Knmi.ref, the daily differences are mainly a result of natural fluctuations. The unventilated Stevenson screens of the second group show much more variation in the daily values than the screens in the first group and their values display positive skewness. The Young.aspII that constitutes the third group shows the largest variation in the daily values with a large negative bias and negative skewness of  $\Delta T_x$  in all seasons. These findings are also reflected in the diurnal cycle differences for winter and summer for four combinations of windspeed and cloudiness.

It is shown that for the calculation of  $T_n$  and  $T_x$  from the 15-s samples the length of the measurements interval  $L$  may be an important factor. Slow responding screens, like the natural ventilated Stevenson screens, have less extreme  $T_n$  and  $T_x$  values than fast-responding screens. The reason is the smaller variation of the 15-s samples in the sliding  $L$ -min intervals for the slow-responding screens compared to the fast-responding screens. For  $T_x$  this effect is a factor 3–4 larger than for  $T_n$ . In practice, when changing from an old screen to a new screen, this effect will often partly compensate for other effects. For instance, a change from Stev.pvc to Knmi.ref or Young.aspII may result in a lowering of  $T_x$  because Knmi.ref and Young.aspII are less sensitive to radiation errors than Stev.pvc. The above-mentioned effect partly compensates for this lowering (depending on the magnitude of  $L$ ). The compensation effect may be even larger when the old screen is supplied with a mercury thermometer and the new screen with a fast-responding sensor like the platinum resistance sensor (e.g. the Pt500). More research is needed to estimate the magnitude of these effects in that case.

We constructed two transfer functions that transfer Knmi.ref into Stev.pvc, one for the 10-min mean temperature values and one for the daily  $T_n$  and  $T_x$ . We have shown that it is feasible to derive such functions, although it was felt that the lack of adequate windspeed data at screen height near the experimental site hampered the derivation.

Windspeed is an important factor in the transfer functions, with its relative importance increasing with decreasing windspeed. Laboratory experiments show that, e.g. the radiation error under conditions of snow cover may increase to about 8°C for windspeeds close to 0 m/s (Gill, 1983). Brandsma (2004) showed the importance of small windspeeds on temperature differences for even slight changes in location of the screen. Measurements of high-accuracy small windspeeds at screen height will therefore be very helpful in developing improved weather-dependent transfer functions in case of future inhomogeneities. The extent of this deserves further

study. In addition, long records of such windspeeds may allow for objectively monitoring the local roughness of temperature sites. We therefore recommend that temperature stations of climatological interest will be equipped with anemometers for high-accuracy (especially in the range 0–3 m/s) wind speed measurements, placed near the thermometers' screen at screen level. These measurements could be performed in addition to the current operational windspeed measurements (mostly at 10-m height as recommended by the WMO).

In contrast to the practice for homogenizing series using mean monthly or mean annual correction, weather-dependent daily or sub-daily corrections using transfer functions will almost always reduce the values of the new measurements to the old. The reason is that the predictors of interest (like global radiation or windspeed at screen level) are mostly not available in the past. On the other hand, using these variables in transfer equations necessitates that continuity of their measurements can be guaranteed towards the future.

The duration that an intercomparison should take depends on the application. For instance, when using a transfer function for 10-min temperatures like the one in Equation (1), 1 year of simultaneous measurements may be long enough to obtain statistically significant coefficients. For the daily  $T_n$  and  $T_x$  in Equations (3) and (4) a period of 2 years may be hardly sufficient. On the other hand, without considering the weather dependence of temperature differences, much longer periods than 2 years may be needed to obtain, e.g. representative monthly mean corrections. In the latter case, it should be kept in mind that the longer the duration of an intercomparison, the more difficult it will be to keep conditions constant (responsible people, environment, measuring methods, etc.).

If adequate modeling of rare events (like periods with snow cover in the Netherlands) is important, it may be advisable to perform part of the screen intercomparison in countries with a comparable climate but with a much more frequent occurrence of the event considered.

The present work focused on the effects of changes in thermometer screens on temperature measurements. The same approach applies to other variables and to other changes in the measurement infrastructure or changes in the local environment surrounding the measurement site. However, the national agencies responsible for carrying out climate observations are not yet equipped for deriving and applying weather-dependent corrections using transfer functions in these cases. Still, when we want to use the time series of the corresponding meteorological variables for studying long-term trends in, e.g. daily or sub-daily extremes, we have to deal with these problems. Because of the labor-intensive character of the related work, it may only be possible for a selected set of high-quality climate monitoring stations (e.g. a subset taken from the GCOS Surface Network; Peterson *et al.*, 1997). We hope that the present work (Parts I and II) will stimulate both the

performance of parallel measurements and the accompanying development and application of transfer functions for these stations.

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